## We examine the asteroseismic phenomenology of stellar rotation when the direction of the rotational axis varies with radial position.

## Motivation

In the shellular approximation of rotation, the angular momenta of all mass shells are assumed to be aligned. However, this may not necessarily be true, as, e.g.:

- Tidal realignment of oblique planets exerts torque on CZ (e.g. Nick Saunders's poster)
- Engulfment deposits angular momentum directly into stellar envelope

Both result in stratification of the axis of rotation!

- What are the observational signatures of this misalignment?
- How will these affect our existing rotational measurements?


## Generalised rotational splitting matrices

Neglecting latitudinal differential rotation, conventional rotational multiplet components are the eigensystem of a linear operator with matrix elements

$$
\left\langle\boldsymbol{\xi}_{\ell m}, \hat{\mathcal{R}} \boldsymbol{\xi}_{\ell m^{\prime}}\right\rangle=\beta_{n \ell} \int \mathbf{J}_{z} \Omega(r) K_{n \ell}(r) \mathrm{d} r, \quad \mathbf{J}_{z}=\operatorname{diag}(-\ell, \ell+1, \ldots, \ell-1, \ell),
$$

in the basis of spherical harmonics.
$\mathbf{J}_{z}$ is one of three basis matrices in the $\ell^{\text {th }}$ representation of $\mathfrak{s o}(3)-$ i.e. of rank $2 \ell+1-$ along with $\mathbf{J}_{x}$ and $\mathbf{J}_{y}$. These satisfy commutation relations

$$
\begin{equation*}
\left[\mathbf{J}_{i}, \mathbf{J}_{j}\right]=-i \epsilon_{i j j} \mathbf{J}_{k} . \tag{2}
\end{equation*}
$$

If at each radial position $r$ we have not only $\Omega(r)$ but in general $\boldsymbol{\Omega}(r)=\Omega(r) \hat{\mathbf{n}}(r)$, then we now need to find the eigensystem of

$$
\mathbf{R}_{n l}=\beta_{n \ell} \int K_{n \ell}(r)[\boldsymbol{\Omega}(r) \cdot \overrightarrow{\mathbf{J}}] \mathrm{d} r .
$$

Eigenvalues of this matrix are generalised multiplet splittings, and eigenvectors - forming some unitary matrix $\mathbf{U} \in \rho_{\ell}[S O(3)]$ - specify a change of basis (in spherical harmonics) from the reference coordinate system, to those of the normal modes. In particular, $\hat{\mathbf{n}} \cdot \overrightarrow{\mathbf{J}}=\mathbf{d}^{\ell} \mathbf{J}_{\mathbf{z}} \mathbf{d}^{\ell}, \dagger$ where $\mathbf{d}^{\ell}$ is Wigner's $d$-matrix: $\Longrightarrow \delta \omega_{m} \sim m \delta \omega_{m=1}$ even with misalignment.

## Two-Zone Model of Misalignment

Consider a rotating core (with constant $\Omega_{\text {core }}$ ) aligned with the $z$-axis, and an envelope (constant $\Omega_{\text {env }}$ ) aligned with unit vector $\hat{n}$; the boundary between the two sits at radius $r_{0}$. We define a relative sensitivity parameter $\alpha$ for a given multiplet, and misalignment angle $\beta$, as

$$
\alpha_{n \ell}=\int_{0}^{r_{0}} K_{n \ell}(r) \mathrm{d} r, \quad \cos \beta=\mathbf{e}_{z} \cdot \hat{\mathbf{n}} .
$$

Effective rotational splitting is then

$$
\delta \omega_{\mathrm{rot}, n \ell} \sim m \beta_{n \ell}\left|\Omega_{\mathrm{eff}}\right|,
$$

where $\Omega_{\text {eff }}$ is found from vector addition (fig. 1). By the triangle inequality, we have that $\left|\boldsymbol{\Omega}_{\text {eff }} \leq \alpha\right| \boldsymbol{\Omega}_{\text {corel }}|+(1-\alpha)| \Omega_{\text {env }} \mid$.

- Misalignment systematically reduces rotational splittings compared to the aligned case!
- Amount depends on both misalignment angle $\beta$ and sensitivity $\alpha$ (fig. 2).


## (5)

## $\alpha \Omega_{\text {core }}$

(6)

Figure 1: Geometric picture of misaligned rotational splittings as addition of angular momentum vectors



Figure 2: Suppression of rotational splitting in presence of misalignment. Left: Schematic representation of rotational spliting when $\Omega_{\text {corr }}=\Omega_{\text {env }}$ with cos $\beta=0$, showing dependence on $\alpha$. Intermediate values of $\alpha$ between 0 and 1 reduce the splitting width. Right: The amount of reduction depends on the value of $\beta$.
The apparent inclination implied by the relative amplitude of multiplet components will also depend on $\alpha$ and $\beta$ (fig. 3).



Figure 3: Dependence of the apparent inclination angle on $\alpha$ and $\beta$. Left: Numerical dependence on $\alpha$ and $\beta$ of an effective $\beta$ (as in fig. 1) associated with the multiplet. Right: Relationship between $i_{\text {corer }} i_{\text {enven }}$ and $i_{\text {eff }}$ on a spherical triangle. Note the geometrical degeneracy owing to an underspecified projection angle $\gamma$.

## Mode Coupling in the Two-Zone Model

If these zones coincide with propagation regions of a mixed-mode solar-like oscillator, rotation interacts with mode coupling via a Quadratic Hermitian Eigenvalue Problem. The unitary matrices $\mathbf{d}^{\ell}$ modify this coupling as
$\left(\left[\begin{array}{cc}-\boldsymbol{\Omega}_{p}^{2} & \mathbf{A} \\ \mathbf{A}^{\dagger} & -\boldsymbol{\Omega}_{g}^{2}\end{array}\right] \otimes \mathbb{I}_{2 \ell+1}+\omega\left[\begin{array}{cc}\mathbf{R}_{p} \otimes \mathbf{d}^{\ell} \mathbf{J}_{z} \mathbf{d}^{\ell, \dagger} & 0 \\ 0 & \mathbf{R}_{g} \otimes \mathbf{J}_{z}\end{array}\right]+\omega^{2}\left[\begin{array}{cc}\mathbb{I}_{p} & \mathbf{D} \\ \mathbf{D}^{\dagger} & \mathbb{I}_{g}\end{array}\right] \otimes \mathbb{I}_{2 \ell+1}\right) \mathbf{c}=0$,
where $\otimes$ is the tensor (Kronecker) product.

- Upon misalignment, mixed modes become linear combinations of $p$ - and g-modes with different $m_{p}$ and $m_{g}$.
- Since $\alpha=\frac{\zeta}{2-\zeta}$, different mixing fractions $\zeta$ will yield different apparent inclinations!
- Suggests reinterpretation of existing inclination measurements made using mixed modes (e.g. Kepler-56: Huber et al., 2013)


## Near-degeneracy Effects: Modified Avoided Crossings <br> Modes of $m_{p} \neq m_{g}$ may couple to each other, depending on the amount of misalignment. The

 morphology of rotational avoided crossings is likewise modified (fig. 4).



$\Omega_{\text {core }} / 2 \pi \mu \mathrm{~Hz}$
Figure 4: Rotational avoided crossings in the presence of rotational misalignment, computed with respect to a red giant MESA model (Model I of Ong et al., 2022). Each set of two panels shows the same set of avoided crossings, coloured variously by either $m_{p}$ or $m_{g}$, for a single value of the core-envelope misalignment angle $\beta$.

## Key Takeaways

- Stratified misalignment systematically reduces rotational splittings, modifies the shape of rotational avoided crossings, and may be diagnosed by per-multiplet variability of the implied inclination angle. However, inferences of the misalignment angle are geometrically degenerate without further constraints from e.g. Rossiter-McLaughlin measurements.


## References

