

Asteroseismic Diagnostics of Misaligned Cores and Envelopes J. M. Joel Ong^{1 2 *}



(1)

(4)

Motivation

In the **shellular approximation** of rotation, the angular momenta of all mass shells are **assumed to be aligned.** However, this may not necessarily be true, as, e.g.:

- **Tidal realignment** of oblique planets exerts **torque** on CZ (e.g. Nick Saunders's poster)
- **Engulfment** deposits angular momentum **directly** into stellar envelope

Both result in stratification of the axis of rotation!

► What are the **observational signatures** of this misalignment?

Mode Coupling in the Two-Zone Model

If these zones coincide with propagation regions of a **mixed-mode solar-like oscillator**, rotation interacts with mode coupling via a **Quadratic Hermitian Eigenvalue Problem**. The unitary matrices \mathbf{d}^{ℓ} modify this coupling as

$$\left(\begin{bmatrix} -\boldsymbol{\Omega}_{p}^{2} & \mathbf{A} \\ \mathbf{A}^{\dagger} & -\boldsymbol{\Omega}_{g}^{2} \end{bmatrix} \otimes \mathbb{I}_{2\ell+1} + \omega \begin{bmatrix} \mathbf{R}_{p} \otimes \mathbf{d}^{\ell} \mathbf{J}_{z} \mathbf{d}^{\ell,\dagger} & 0 \\ 0 & \mathbf{R}_{g} \otimes \mathbf{J}_{z} \end{bmatrix} + \omega^{2} \begin{bmatrix} \mathbb{I}_{p} & \mathbf{D} \\ \mathbf{D}^{\dagger} & \mathbb{I}_{g} \end{bmatrix} \otimes \mathbb{I}_{2\ell+1} \right) \mathbf{c} = 0, \quad (7)$$

► How will these affect our **existing rotational measurements**?

Generalised rotational splitting matrices

Neglecting latitudinal differential rotation, conventional rotational multiplet components are the **eigensystem of a linear operator** with matrix elements

$$\left\langle \boldsymbol{\xi}_{\ell m}, \hat{\mathcal{R}} \boldsymbol{\xi}_{\ell m'} \right\rangle = \beta_{n\ell} \int \mathbf{J}_z \Omega(r) K_{n\ell}(r) \, \mathrm{d}r, \ \mathbf{J}_z = \operatorname{diag}\left(-\ell, \ell+1, \dots, \ell-1, \ell\right),$$

in the basis of spherical harmonics.

 \mathbf{J}_z is one of three basis matrices in the ℓ^{th} representation of $\mathfrak{so}(3)$ — i.e. of rank $2\ell + 1$ — along with \mathbf{J}_x and \mathbf{J}_y . These satisfy commutation relations

$$[\mathbf{J}_i, \mathbf{J}_j] = -i\epsilon_{ijk}\mathbf{J}_k.$$
⁽²⁾

If at each radial position r we have not only $\Omega(r)$ but in general $\Omega(r) = \Omega(r)\hat{\mathbf{n}}(r)$, then we now need to find the eigensystem of

$$\mathbf{R}_{nl} = \beta_{n\ell} \int K_{n\ell}(r) \left[\mathbf{\Omega}(r) \cdot \vec{\mathbf{J}} \right] \, \mathrm{d}r.$$
(3)

 $\alpha \Omega_{\mathsf{COre}}$

Eigenvalues of this matrix are **generalised multiplet splittings**, and eigenvectors — forming some unitary matrix $\mathbf{U} \in \rho_{\ell}[SO(3)]$ — specify a change of basis (in spherical harmonics) from the reference coordinate system, to those of the normal modes. In particular, $\hat{\mathbf{n}} \cdot \vec{\mathbf{J}} = \mathbf{d}^{\ell} \mathbf{J}_z \mathbf{d}^{\ell,\dagger}$, where \mathbf{d}^{ℓ} is Wigner's *d*-matrix: $\implies \delta \omega_m \sim m \delta \omega_{m=1}$ even with misalignment.

Two-Zone Model of Misalignment

Consider a rotating core (with constant Ω_{core}) aligned with the *z*-axis, and an envelope (constant Ω_{env}) aligned with unit vector $\hat{\mathbf{n}}$; the boundary between the two sits at radius r_0 . We define a relative sensitivity parameter α for a given multiplet, and misalignment angle β , as

- where \otimes is the tensor (Kronecker) product.
- Upon misalignment, mixed modes become linear combinations of p- and g-modes with different m_p and m_q .
- Since $\alpha = \frac{\zeta}{2-\zeta}$, different mixing fractions ζ will yield different apparent inclinations!
- Suggests reinterpretation of existing inclination measurements made using mixed modes (e.g. Kepler-56: Huber et al., 2013)

Near-degeneracy Effects: Modified Avoided Crossings

Modes of $m_p \neq m_q$ may couple to each other, depending on the amount of misalignment. The morphology of rotational avoided crossings is likewise modified (fig. 4).



$$\alpha_{n\ell} = \int_0^{r_0} K_{n\ell}(r) \, \mathrm{d}r; \quad \cos\beta = \mathbf{e}_z \cdot \hat{\mathbf{n}}.$$

(5)

(6)

Effective rotational splitting is then

$$\delta \omega_{\mathrm{rot},n\ell} \sim m \beta_{n\ell} |\mathbf{\Omega}_{\mathrm{eff}}|,$$

where $\Omega_{
m eff}$ is found from vector addition (fig. 1). By the triangle inequality, we have that

$$|\mathbf{\Omega}_{\mathsf{eff}}| \le \alpha |\mathbf{\Omega}_{\mathsf{core}}| + (1-\alpha) |\mathbf{\Omega}_{\mathsf{env}}|.$$

- Misalignment systematically reduces rotational splittings compared to the aligned case!
- Amount depends on **both** misalignment angle β and sensitivity α (fig. 2).





 $\Omega_{ ext{eff}}$

Figure 1: Geometric picture of misaligned

rotational splittings as addition of angular



The **apparent inclination** implied by the relative amplitude of multiplet components will **also** depend on α and β (fig. 3).



Figure 3: Dependence of the apparent inclination angle on α and β . Left: Numerical dependence on α and β of an effective β (as in fig. 1) associated with the multiplet. **Right**: Relationship between i_{core} , i_{env} , and i_{eff} , on a spherical triangle. Note the geometrical degeneracy owing to an underspecified projection angle γ .

Figure 4: Rotational avoided crossings in the presence of rotational misalignment, computed with respect to a red giant MESA model (Model I of Ong et al., 2022). Each set of two panels shows the same set of avoided crossings, coloured variously by either m_p or m_q , for a single value of the core-envelope misalignment angle β .

Key Takeaways

Stratified misalignment systematically reduces rotational splittings, modifies the shape of rotational avoided crossings, and **may be diagnosed** by per-multiplet variability of the implied inclination angle. However, inferences of the misalignment angle are **geometrically degenerate** without further constraints from e.g. Rossiter-McLaughlin measurements.

References

Huber, D., Carter, J. A., Barbieri, M., et al. 2013, Science, 342, 331 Ong, J. M. J., Bugnet, L., & Basu, S. 2022, ApJ, 940, 18